

# Numerical Solutions to Differential Equations

## Lecture Notes — The Finite Element Method #2

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## Outline

- 1 **The Finite Element Method**
  - Recap
  - Looking for Solutions...
- 2 **Stiffness Matrix, and Load Vector**
  - Identifying a Linear System...
  - Properties of the Stiffness Matrix
- 3 **Error Estimation**
  - Building the Toolbox
  - Error Control
- 4 **FEM for Partial Differential Equations**
  - Initial Example: The Poisson Equation
  - Analytical Tools: Vector Calculus
  - Variational Formulation
  - Minimization Formulation
  - The Road Ahead...

## Quick Recap

Last time we formulated the Galerkin (variational,  $(V_h)$ ) and Ritz (minimization,  $(M_h)$ ) methods which will give us the finite element solutions to the differential equation  $(D)$ :

$$-u'' = f, \quad + \quad \text{Dirichlet Boundary Conditions}$$

**The Function Space  $V_h$ :** Given the basis functions  $\phi_i$  (the “tent” functions) we can write

$$V_h = \left\{ v : v = \sum_{i=1}^N \eta_i \phi_i(x) \right\}.$$

## Quick Recap

**Inner Products:** We defined the two (for now identical) inner products:

$$a(u, v) = \int_{\mathcal{I}} uv \, dx, \quad (u, v) = \int_{\mathcal{I}} uv \, dx.$$

Further, we defined the energy **functional**:

$$F(u) = \frac{1}{2}a(v', v') - (f, v).$$

Given these definitions we formulated the Galerkin and Ritz problems...

## Quick Recap



### Galerkin's Method (Variational Approach)

$$(V_h) \quad \text{Find } u_h \in V_h \text{ so that } a(u_h', v_h') = (f, v_h) \quad \forall v_h \in V_h.$$

### Ritz' Method (Minimization Approach)

$$(M_h) \quad \text{Find } u_h \in V_h \text{ so that } F(u_h) \leq F(v_h) \quad \forall v_h \in V_h.$$

## Looking for the Solution...

If  $u_h \in V_h$  is a solution to  $(V_h)$ , then **in particular**

$$a(u_h', \phi_j') = (f, \phi_j), \quad j = 1, 2, \dots, N.$$

Also, we can write  $u_h$  in terms of the basis functions:

$$u_h(x) = \sum_{j=1}^N \xi_j \phi_j(x), \quad \xi_j = u_h(x_j).$$

Therefore we can rewrite the above equation

$$\sum_{i=1}^N \xi_i a(\phi_i', \phi_j') = (f, \phi_j), \quad j = 1, 2, \dots, N.$$

## The Stiffness Matrix and Load Vector

From

$$\sum_{i=1}^N \xi_i a(\phi'_i, \phi'_j) = (f, \phi_j), \quad j = 1, 2, \dots, N$$

we identify the vectors  $\tilde{\xi} = \{\xi_1, \xi_2, \dots, \xi_N\}^T$ , and  $\tilde{\mathbf{b}} = \{b_1, b_2, \dots, b_N\}^T$ , where  $b_i = (f, \phi_i)$ , and the matrix  $A$ , where  $A_{ij} = a(\phi'_i, \phi'_j)$ .

We can now write our problem as

$$A\tilde{\xi} = \tilde{\mathbf{b}}.$$

For historical reasons (Structural Mechanics)  $A$  is known as the **stiffness matrix**, and  $\tilde{\mathbf{b}}$  as the **load vector**.

## Computing the Stiffness Matrix

The elements of the stiffness matrix can be computed: —

First we notice that if the basis is the piecewise linear “tent-functions”, then if  $|i - j| > 1$ , then  $\phi_i$  and  $\phi_j$  are non-overlapping which means  $a(\phi'_i, \phi'_j) = 0$ .

**The Diagonal Entries**,  $j = 1, 2, \dots, N$ .

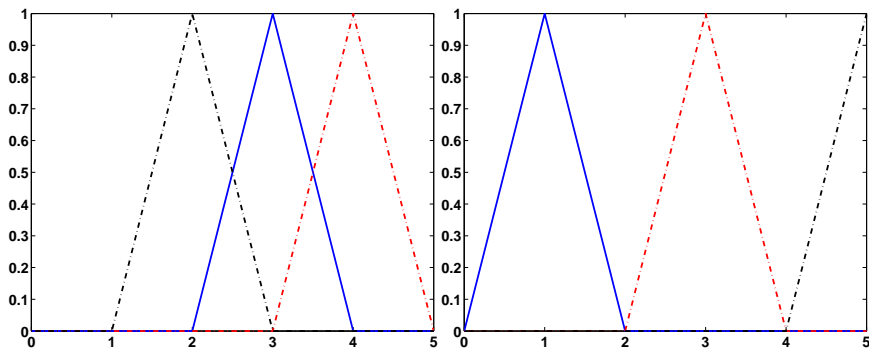
$$a(\phi'_j, \phi'_j) = \int_{x_{j-1}}^{x_j} \frac{1}{h_j^2} dx + \int_{x_j}^{x_{j+1}} \frac{1}{h_{j+1}^2} dx = \frac{1}{h_j} + \frac{1}{h_{j+1}}$$

**The Super- and Sub-Diagonal Entries**,  $j = 2, 3, \dots, N$ .

$$a(\phi'_{j-1}, \phi'_j) = a(\phi'_j, \phi'_{j-1}) = \int_{x_{j-1}}^{x_j} \frac{-1}{h_j} \frac{1}{h_j} dx = -\frac{1}{h_j}$$



## Illustration: Overlapping and Nonoverlapping Tent-Functions



## The Stiffness Matrix is Symmetric Positive Definite (SPD)

$A$  is **symmetric** since  $a(u, v) = a(v, u)$ , and with  $v(x) = \sum_{i=1}^N \eta_j \phi_i(x)$  we get

$$\sum_{i,j=1}^N \eta_i a(\phi'_i, \phi'_j) \eta_j = a \left( \sum_{i=1}^N \eta_i \phi'_i, \sum_{j=1}^N \eta_j \phi'_j \right) = \int_1 [v'_h(x)]^2 dx \geq 0$$

Equality holds if and only if  $v'(x) \equiv 0$ , which implies  $v(x) \equiv 0$  by the boundary conditions.

**Fact:** A matrix  $A$  is SPD if

$$x^T A x > 0, \quad \forall x \in \mathbb{R}^n \setminus \{0\}$$

**Fact:** An SPD matrix (i) has positive eigenvalues, (ii) is non-singular.

## The Special Case $h_j = h, j = 1, 2, \dots, N$

If we equi-partition the interval we get the linear system

$$\frac{1}{h} \begin{bmatrix} 2 & -1 & 0 & \dots & \dots & 0 \\ -1 & 2 & -1 & \ddots & & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & & \ddots & -1 & 2 & -1 \\ 0 & \dots & \dots & 0 & -1 & 2 \end{bmatrix} \xi = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ \vdots \\ \vdots \\ b_N \end{bmatrix}$$

If we divide by  $h$  the we recover the standard **finite difference method** for the problem, where the right hand-side

$$b_j = \int_{x_{j-1}}^{x_{j+1}} f(x) \phi_j dx$$

is a **weighted mean** of  $f(x)$  over the interval  $[x_{j-1}, x_{j+1}]$ .

## Estimating the Error for the Model Problem

We are now going to look at the error  $(\mathbf{u} - \mathbf{u}_h)$  where  $u$  is the exact solution of the differential equation, and  $u_h$  is the solution to the finite element problem ( $V_h$ ).

Since  $V_h \subset V$ ,

$$\begin{aligned} a(u', v_h') &= (f, v_h), & \forall v_h \in V_h \\ a(u_h', v_h') &= (f, v_h), & \forall v_h \in V_h \\ \hline a((u - u_h)', v_h') &= 0, & \forall v_h \in V_h \end{aligned}$$

That means that the error is orthogonal to the function space  $V_h$  (as measured by  $a(o', o')$ .)

## Definition and Tools

Definition (The  $L_2$ -norm ( $L_{2,a}$ ?))

$$\|w\|_2 = \sqrt{a(w, w)} = \sqrt{\int_0^1 w^2 dx}$$

Theorem (Cauchy's Inequality)

$$|a(v, w)| \leq \|v\| \|w\|$$

## Cauchy's Inequality: Proof

### Cauchy's Inequality.

Given  $v$  and  $w$ , define the renormalized functions  $\hat{v} = \frac{v}{\|v\|}$  and  $\hat{w} = \frac{w}{\|w\|}$ .  
Which means  $\|\hat{v}\| = \|\hat{w}\| = 1$ . Now

$$\begin{aligned} 0 &\leq \|\pm \hat{v} - \hat{w}\|^2 = a(\pm \hat{v} - \hat{w}, \pm \hat{v} - \hat{w}) \\ &= a(\hat{v}, \hat{v}) \mp 2a(\hat{v}, \hat{w}) + a(\hat{w}, \hat{w}) \\ &= 2 \mp 2a(\hat{v}, \hat{w}). \end{aligned}$$

Hence,

$$|a(\hat{v}, \hat{w})| \leq 1.$$

Removing the linear normalization factors give

$$|a(v, w)| \leq \|v\| \|w\|.$$



## Error Control — Theorem

### Theorem

For any  $v_h \in V_h$  we have

$$\|(u - u_h)'\| \leq \|(u - v_h)'\|.$$

So, measured in the  $L_2$ -norm of the derivative (also known as the  $H^1$ -seminorm), the solution  $u_h$  to the discrete problem is closer to the solution  $u$  to the original continuous ODE than any other function in the function space  $V_h$ .

## Error Control — Proof

### Proof.

Let  $v_h \in V_h$  be arbitrary, and set  $w_h = u_h - v_h$ . Then  $w_h \in V_h$  and we get:

$$\begin{aligned} \|(u - u_h)'\|^2 &= a((u - u_h)', (u - u_h)') + \underbrace{a((u - u_h)', w_h')}_{=0} \\ &= a((u - u_h)', (u - u_h + w_h)') \\ &= a((u - u_h)', (u - v_h)') \\ \{\text{Cauchy's}\} &\leq \|(u - u_h)'\| \|(u - v_h)'\|. \end{aligned}$$

Dividing through by  $\|(u - u_h)'\|$  gives

$$\|(u - u_h)'\| \leq \|(u - v_h)'\|.$$





## Error Control — Applying the Theorem

From the theorem we can get a quantitative estimate (upper bound) for the error  $\|(u - u_h)'\|$  by estimating  $\|(u - v_h)'\|$  where  $v_h \in V_h$  is a suitably chosen function.

Let  $v_h$  be the **linear interpolant** of  $u$ , i.e.  $v_h(x_j) = u(x_j)$ , then

$$|u'(x) - v_h'(x)| \leq h \max_{x \in [0,1]} |u''(x)| \quad [1]$$

$$|u(x) - v_h(x)| \leq \frac{h^2}{8} \max_{x \in [0,1]} |u''(x)| \quad [2]$$

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$$|u(x) - v_h(x)| \leq \frac{h^2}{8} \max_{x \in [0,1]} |u''(x)| \quad [2]$$

The first result and the theorem gives

$$\|(u - u_h)'\| \leq C_1 h \max_{x \in [0,1]} |u''(x)|.$$

Similarly, from [2], we can get

$$\|u(x) - u_h(x)\| \leq C_2 h^2 \max_{x \in [0,1]} |u''(x)|.$$

## FEM for PDEs — the Poisson Equation

Consider the following boundary value problem

$$\begin{aligned} -\Delta u &= f && \text{in } \Omega \\ u &= 0 && \text{on } \Gamma = \partial\Omega, \end{aligned}$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^2 = \{(x_1, x_2) : x_1 \in \mathbb{R}, x_2 \in \mathbb{R}\}$ .  
 $\Gamma$  is the boundary of  $\Omega$ ,  $f(x_1, x_2)$  a given function, and

$$\Delta u = \left[ \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} \right] u(x_1, x_2) = u_{x_1 x_1}(x_1, x_2) + u_{x_2 x_2}(x_1, x_2).$$

**Physical problems:** Our simple 1-D model problems from last lecture carry over to this 2-D model: heat distribution in a plate, the displacement of an elastic membrane fixed at the boundary under transverse load, etc...

## Vector Calculus — Review (or Crash Course)

We need to generalize “integration by parts” to higher dimensions... We start with the **divergence theorem**:

$$\int_{\Omega} \nabla \circ \tilde{\mathbf{u}} \, d\tilde{\mathbf{x}} = \int_{\Gamma} \tilde{\mathbf{u}} \circ \tilde{\mathbf{n}} \, ds,$$

where  $\tilde{\mathbf{u}} = \{u_1, u_2\}^T$ , and

$$\nabla \circ \tilde{\mathbf{u}} = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2}.$$

$\tilde{\mathbf{n}} = \{n_1, n_2\}^T$  is the unit (length-one, outward) normal to  $\Gamma$ .  $d\tilde{\mathbf{x}}$  denotes area-integration, and  $ds$  integration along the boundary. *c.f.* in 1-D:

$$\int_{[a,b]} \frac{du}{dx} \, dx = \int_{\{a,b\}} u(\mp 1) \, ds = -u(a) + u(b)$$

## Vector Calculus — Review (or Crash Course)

If we apply the divergence theorem to the vector functions  $\tilde{\mathbf{u}}_1 = (vw, 0)$  and  $\tilde{\mathbf{u}}_2 = (0, vw)$ , we find

$$\int_{\Omega} \frac{\partial v}{\partial x_i} w + v \frac{\partial w}{\partial x_i} d\tilde{\mathbf{x}} = \int_{\Gamma} v w n_i ds, \quad i = 1, 2.$$

Let  $\nabla v$  denote the gradient

$$\nabla v = \left\{ \frac{\partial v}{\partial x_1}, \frac{\partial v}{\partial x_2} \right\}^T,$$

and use the result above to write **Green's Formula**

$$\begin{aligned} \int_{\Omega} \nabla v \circ \nabla w d\tilde{\mathbf{x}} &\equiv \int_{\Omega} \left[ \frac{\partial v}{\partial x_1} \frac{\partial w}{\partial x_1} + \frac{\partial v}{\partial x_2} \frac{\partial w}{\partial x_2} \right] d\tilde{\mathbf{x}} \\ &= \int_{\Gamma} \left[ v \frac{\partial w}{\partial x_1} n_1 + v \frac{\partial w}{\partial x_2} n_2 \right] ds - \int_{\Omega} v \left[ \frac{\partial^2 w}{\partial x_1^2} + \frac{\partial^2 w}{\partial x_2^2} \right] d\tilde{\mathbf{x}} \\ &= \int_{\Gamma} v \frac{\partial w}{\partial n} ds - \int_{\Omega} v \Delta w d\tilde{\mathbf{x}}. \end{aligned}$$

## Vector Calculus — Review (or Crash Course)

The normal derivative

$$\frac{\partial w}{\partial n} = \frac{\partial w}{\partial x_1} n_1 + \frac{\partial w}{\partial x_2} n_2 = \nabla w \circ \tilde{\mathbf{n}}$$

is the derivative in the outward normal direction to the boundary  $\Gamma$ .

With this notation we are ready to state the variational formulation corresponding to the differential equation...

## Variational Formulation

(V) Find  $u \in V$  so that  $a(u, v) = (f, v)$ ,  $\forall v \in V$ .

where

$$a(u, v) = \int_{\Omega} \nabla u \circ \nabla v \, d\tilde{\mathbf{x}}$$

$$(f, v) = \int_{\Omega} f v \, d\tilde{\mathbf{x}}$$

$$V = \left\{ \begin{array}{l} v : v \in C(\Omega) \\ v_{x_1} \text{ and } v_{x_2} \text{ are piecewise continuous in } \Omega \\ v = 0 \text{ on } \Gamma \end{array} \right\}$$

**Note:** We have slightly changed the notation of the  $a(\circ, \circ)$  inner-product. This is very common in the PDE framework. The  $a(\circ, \circ)$  inner-product generically involves integrating the product(s) of derivative(s) of the two arguments over the domain.

## Minimization Formulation

$u$  is a solution to  $(V)$  if and only if it is also a solution to the following minimization problem.

$$(M) \quad \text{Find } u \in V \text{ such that } F(u) \leq F(v) \quad \forall v \in V,$$

where  $F(v)$  is the potential energy

$$F(v) = \frac{1}{2}a(v, v) - (f, v).$$



Next...

Constructing a Finite Element Method for this problem...